Response Function of the BGO and NaI(Tl) Detectors Using Monte Carlo Simulations

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INTRODUCTION

The high efficiency of the BGO detectors makes them a very attractive candidates to replace NaI(Tl) detectors, which are widely used in studies of body composition. In this work, the response functions of the BGO and NaI(Tl) detectors were determined at 0.662, 4.4, and 10.0 MeV using three different Monte Carlo codes; EGS4¹, MCNP², and PHOTON³. These codes differ in their input files and transport calculations, and were used to verify the internal consistency of the set-up and of the input data. The energy range of 0.662-10 MeV was chosen to cover energies of interest in body-composition studies. The superior efficiency of the BGO detectors has to be weighed against their inferior resolution, and their higher price than that of the NaI detectors. Since the price of the BGO detectors strongly depends on the crystal's size, its optimization is an important component in the design of the entire system.

METHOD

The intrinsic response functions were calculated for a point source at 10 cm from the base of a cylindrical detector, where the source emission was forced into the solid angle subtended by the detector. The energy deposition in the detector, i.e., the photon spectrum, was calculated using 512 bins 20 keV wide. Special attention was paid to the energy cut-off used in the codes. While this value is fixed at 10 keV in the PHOTON code, to improve accuracy we altered it in the other two codes to 1 keV. The main effect of the cut-off value is in the valley area beside the photopeak (Fig 1a.). The calculated photon spectra were convoluted with an energy-resolution function of the type $a+bE^{\frac{1}{2}}+cE$ that was fitted to the experimental data for each detector reported in the literature⁴. The convolution consisted of Monte Carlo normal sampling of each energy bin. We found that from about 2 MeV and above, the results from the PHOTON code (Figs. 1a, 1b, and 1c). We attributed this discrepancy to inadequate bremsstrahlung cross-sections used in the PHOTON code (we are verifying this with the authors of the code). The Photon code was not used from further at high energies.

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FIGURE 1. Comparison between the results of MCNP4B2, EGS4, and PHOTON Monte Carlo codes: (a) at a source energy of 0.662 MeV; (b) at 4.4 MeV; and (c) at 10 MeV.



FIGURE 2. Simulated 4.4 MeV pulse-height distribution with the specific energy-resolution function for BGO 5"x 3", BGO 3"x 3", and NaI(Tl) 6"x 6" using the EGS4 Monte Carlo code.



FIGURE 3. Simulated 10 MeV pulse-height distribution with the specific energy-resolution function for BGO 5"x 3", BGO 3"x 3" and NaI(Tl) 6"x 6" using the EGS4 Monte Carlo code.

DISCUSSION AND CONCLUSIONS

At high energies, the higher energy resolution of the BGO detectors had little effect on the distribution of pulse height compared to that of the NaI detectors (Figs. 2 and 3). However, the high efficiency of the BGO detectors led to a higher peak-to-background ratio (Fig. 4). At higher



FIGURE 4. Calculated peak-to-total ratios versus incident photon energy for each detector by type and by size.

	$E = 4.4 M_{\odot} V$					$E = 10 M_{\odot} V$				
	E = 4.4 MeV					E - 10 MeV				
Detector	P/N	SE/N	DE/N	P/SE	P/DE	P/N*	SE/N*	DE/N*	P/SE	P/DE
NaI(Tl)										
3"x3"	0.082	0.051	0.019	1.6	4.2	0.031	0.045	0.016	0.7	1.9
6"x6"	0.218	0.066	0.009	3.3	25.7	0.124	0.080	0.009	1.7	13.4
BGO										
3"x3"	0.354	0.074	0.006	4.8	62.3	0.240	0.078	0.007	3.1	35.1
5"x3"	0.427	0.064	0.004	6.7	112.4	0.323	0.070	0.004	4.6	72.6

TABLE 1. The yield per history of the single-escape (SE) and double-escape (DE) peaks and the photopeak-to-escape-peak ratios at two source energies.

ABBREVIATIONS: $N = 10^6$, Number of histories; $N^* = 10^5$, Number of histories; P = photopeak counts; SE = single escape peak net counts; DE = double escape peak net counts.



FIGURE 5. The BGO detector's intrinsic efficiency (E_f) dependence on volume fitted to the logistical equation (the parameters are listed for three incident-photon energies).

energies the shape of the Compton region is almost independent of the type and the size of the detector. For the BGO detector, the photopeak intrinsic efficiency versus detector volume is effectively represented by a single logistical equation of the form

$$E_{f} = \frac{A1 - A2}{1 - V^{0.33}} + A2$$

as shown in Fig. 5. The main differences in the response function for the detectors of different size and type were in the escape peaks, (see TABLE 1). In summary, the BGO detectors, despite their worse energy resolution, appear to be superior to NaI(Tl) detectors for high-energy gamma detection.

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